

2: Three Different Fourier Transforms

- Fourier Transforms
- Convergence of DTFT
- DTFT Properties
- DFT Properties
- Symmetries
- Parseval's Theorem
- Convolution
- Sampling Process
- Zero-Padding
- Phase Unwrapping
- Uncertainty principle
- Summary
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Forward Transform	Inverse Transform
CTFT	$X(j\Omega) = \int_{-\infty}^{\infty} x(t)e^{-j\Omega t}dt$
	$x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\Omega)e^{j\Omega t}d\Omega$

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DTFT $X(e^{j\omega}) = \sum_{-\infty}^{\infty} x[n]e^{-j\omega n}$	$x[n] = \frac{1}{2\pi} \int_{-\pi}^{\pi} X(e^{j\omega})e^{j\omega n} d\omega$

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We use Ω for “real” and $\omega = \Omega T$ for “normalized” angular frequency.
Nyquist frequency is at $\Omega_{\text{Nyq}} = 2\pi \frac{f_s}{2} = \frac{\pi}{T}$ and $\omega_{\text{Nyq}} = \pi$.

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DFT	$X[k] = \sum_0^{N-1} x[n]e^{-j2\pi \frac{kn}{N}}$	$x[n] = \frac{1}{N} \sum_0^{N-1} X[k]e^{j2\pi \frac{kn}{N}}$

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For “power signals” (energy \propto duration), CTFT & DTFT are unbounded.
Fix this by normalizing:

$$X(j\Omega) = \lim_{A \rightarrow \infty} \frac{1}{2A} \int_{-A}^A x(t)e^{-j\Omega t} dt$$
$$X(e^{j\omega}) = \lim_{A \rightarrow \infty} \frac{1}{2A+1} \sum_{-A}^A x[n]e^{-j\omega n}$$

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Convergence of DTFT

DTFT: $X(e^{j\omega}) = \sum_{-\infty}^{\infty} x[n]e^{-j\omega n}$ does not converge for all $x[n]$.

Consider the finite sum: $X_K(e^{j\omega}) = \sum_{-K}^K x[n]e^{-j\omega n}$

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Strong Convergence:

$x[n]$ absolutely summable $\Rightarrow X(e^{j\omega})$ converges **uniformly**

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Weaker convergence:

$x[n]$ finite energy $\Rightarrow X(e^{j\omega})$ converges in **mean square**

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Weaker convergence:

$x[n]$ finite energy $\Rightarrow X(e^{j\omega})$ converges in **mean square**
 $\sum_{-\infty}^{\infty} |x[n]|^2 < \infty \Rightarrow \frac{1}{2\pi} \int_{-\pi}^{\pi} |X(e^{j\omega}) - X_K(e^{j\omega})|^2 d\omega \xrightarrow[K \rightarrow \infty]{} 0$

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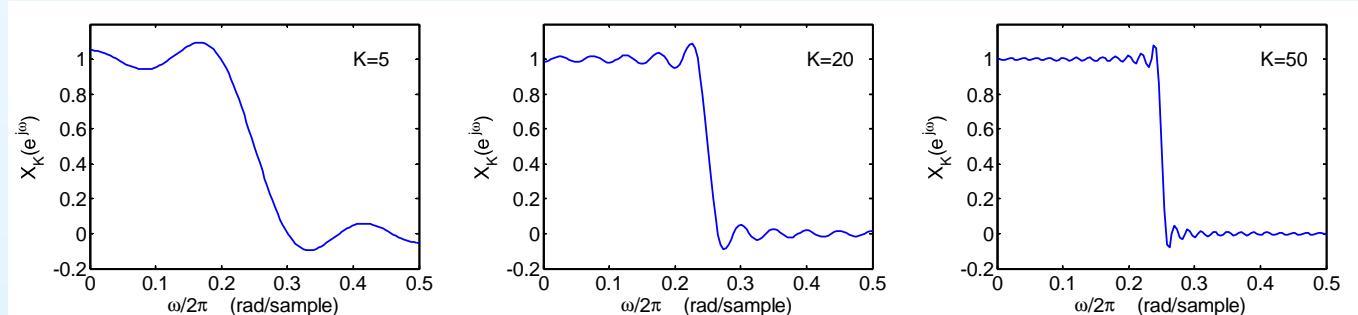
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Example: $x[n] = \frac{\sin 0.5\pi n}{\pi n}$



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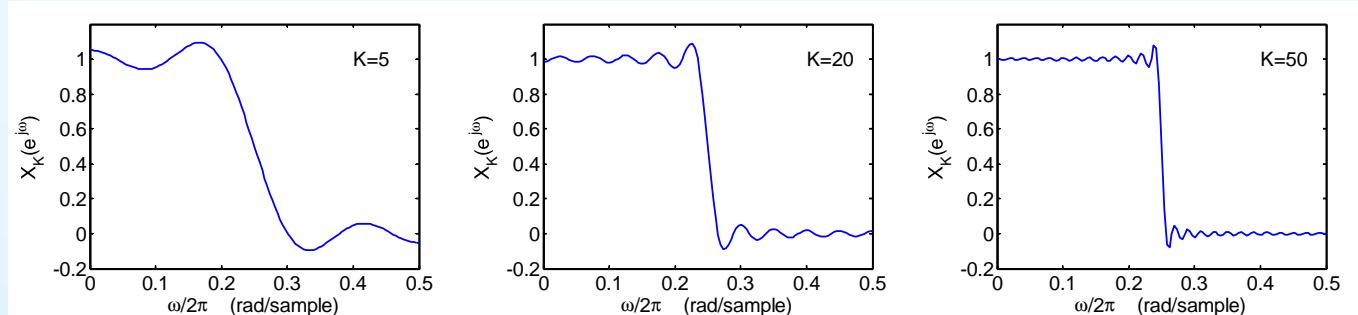
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Gibbs phenomenon:

Converges at each ω as $K \rightarrow \infty$ but peak error does not get smaller.

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- DTFT is **periodic** in ω : $X(e^{j(\omega+2m\pi)}) = X(e^{j\omega})$ for integer m .

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$$\begin{aligned} \text{Proof: } X(e^{j\omega}) &= \sum_{n=-\infty}^{\infty} x[n]e^{-j\omega n} \\ &= \sum_{n=-\infty}^{\infty} x[n] \int_{-\infty}^{\infty} \delta(t - nT)e^{-j\omega \frac{t}{T}} dt \\ &\stackrel{(i)}{=} \int_{-\infty}^{\infty} \sum_{n=-\infty}^{\infty} x[n]\delta(t - nT)e^{-j\omega \frac{t}{T}} dt \end{aligned}$$

(i) OK if $\sum_{n=-\infty}^{\infty} |x[n]| < \infty$.

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Equivalent to multiplying a continuous $x(t)$ by an impulse train.

$$\begin{aligned} \text{Proof: } X(e^{j\omega}) &= \sum_{n=-\infty}^{\infty} x[n]e^{-j\omega n} \\ &= \sum_{n=-\infty}^{\infty} x[n] \int_{-\infty}^{\infty} \delta(t - nT)e^{-j\omega \frac{t}{T}} dt \\ &\stackrel{(i)}{=} \int_{-\infty}^{\infty} \sum_{n=-\infty}^{\infty} x[n]\delta(t - nT)e^{-j\omega \frac{t}{T}} dt \\ &\stackrel{(ii)}{=} \int_{-\infty}^{\infty} x_{\delta}(t)e^{-j\Omega t} dt \end{aligned}$$

(i) OK if $\sum_{n=-\infty}^{\infty} |x[n]| < \infty$. (ii) use $\omega = \Omega T$.

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$$\text{DFT: } X[k] = \sum_0^{N-1} x[n]e^{-j2\pi \frac{kn}{N}}$$

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DFT equals the normalized DTFT

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One domain	Other domain
Discrete	Periodic
Symmetric	Symmetric
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Fourier transforms preserve “energy”

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CTFT $\int |x(t)|^2 dt = \frac{1}{2\pi} \int |X(j\Omega)|^2 d\Omega$

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More generally, they actually preserve **complex inner products**:

$$\sum_0^{N-1} x[n]y^*[n] = \frac{1}{N} \sum_0^{N-1} X[k]Y^*[k]$$

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If we regard \mathbf{x} and \mathbf{X} as vectors, then $\mathbf{X} = \mathbf{F}\mathbf{x}$ where \mathbf{F} is a symmetric matrix defined by $f_{k+1,n+1} = e^{-j2\pi \frac{kn}{N}}$.

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The inverse DFT matrix is $\mathbf{F}^{-1} = \frac{1}{N} \mathbf{F}^H$
equivalently, $\mathbf{G} = \frac{1}{\sqrt{N}} \mathbf{F}$ is a **unitary matrix** with $\mathbf{G}^H \mathbf{G} = \mathbf{I}$.

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DTFT: Convolution → Product

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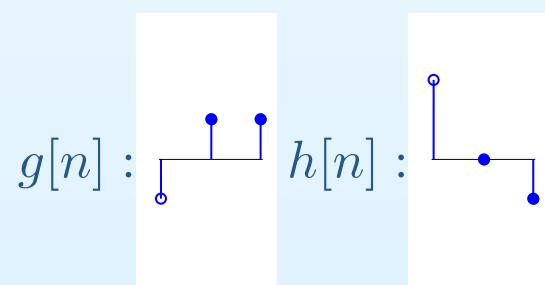
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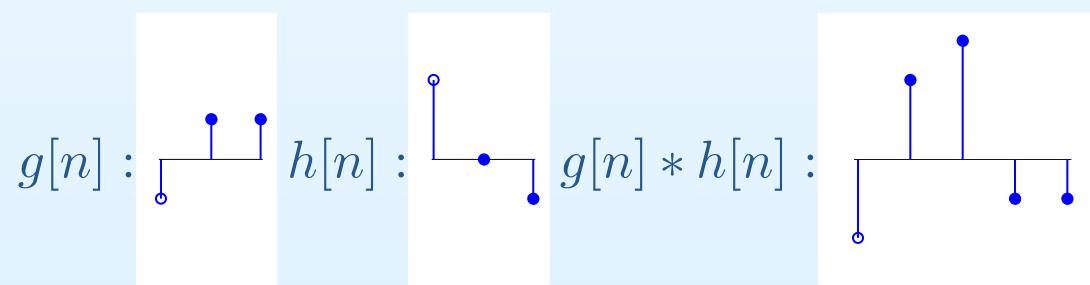
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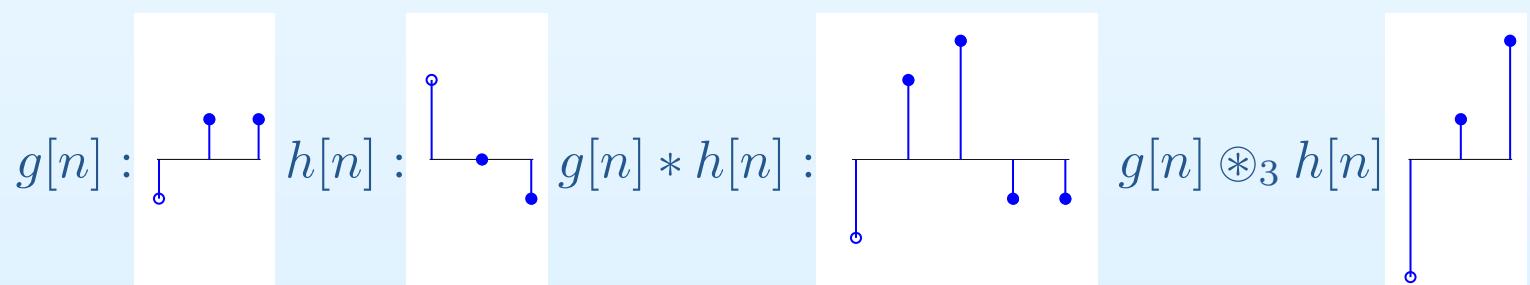
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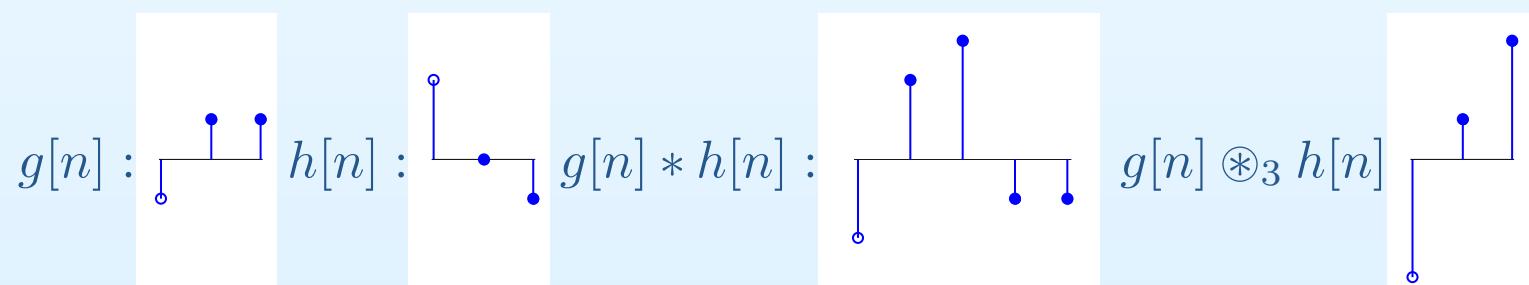
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DTFT: Product → Circular Convolution $\div 2\pi$

$$y[n] = g[n]h[n]$$



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DTFT: Convolution → Product

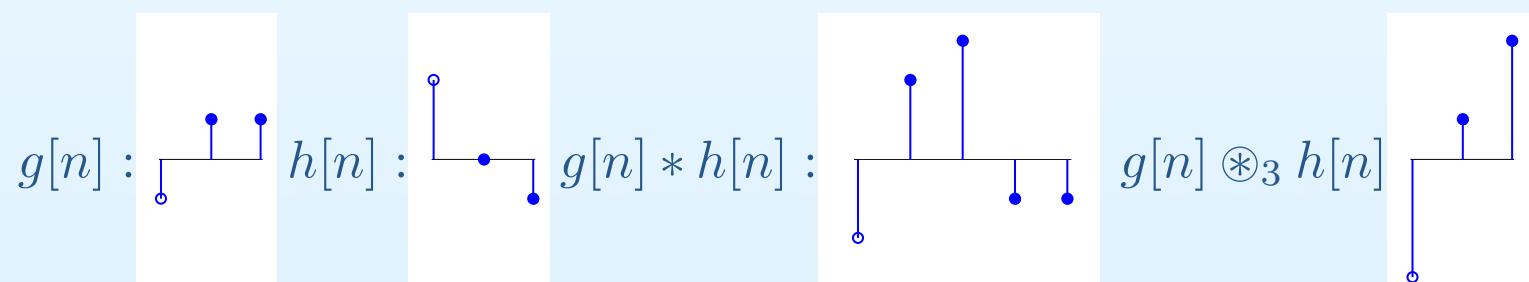
$$x[n] = g[n] * h[n] = \sum_{k=-\infty}^{\infty} g[k]h[n-k]$$
$$\Rightarrow X(e^{j\omega}) = G(e^{j\omega})H(e^{j\omega})$$

DFT: Circular convolution → Product

$$x[n] = g[n] \circledast_N h[n] = \sum_{k=0}^{N-1} g[k]h[(n-k)_{\text{mod } N}]$$
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DTFT: Product → Circular Convolution $\div 2\pi$

$$y[n] = g[n]h[n]$$
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$$\frac{1}{2\pi} \int_{-\pi}^{\pi} G(e^{j\theta})H(e^{j(\omega-\theta)})d\theta$$



Convolution

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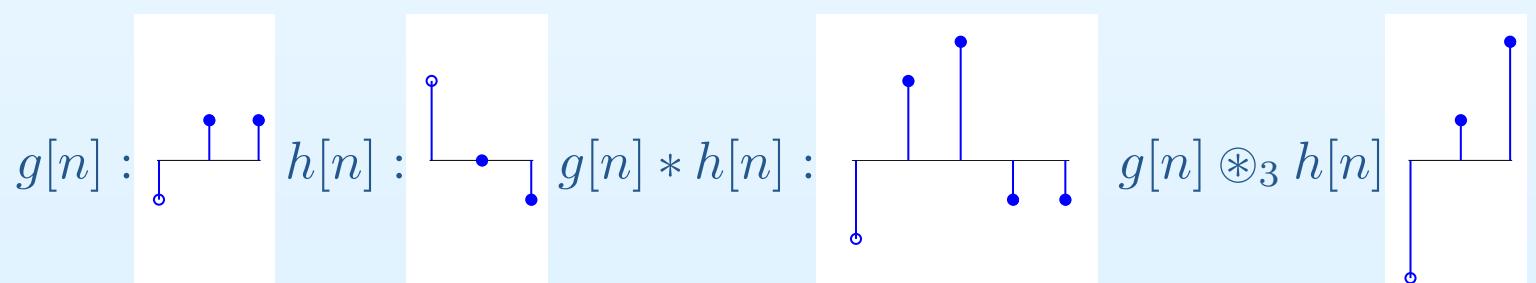
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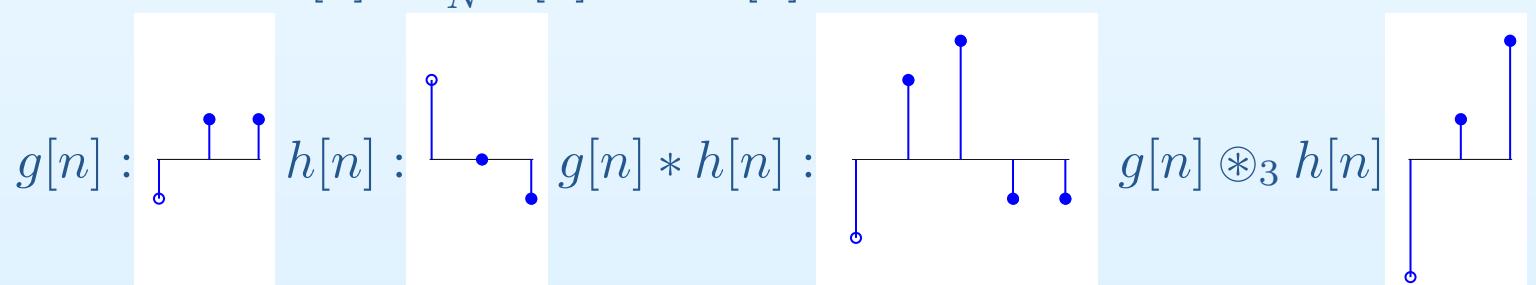
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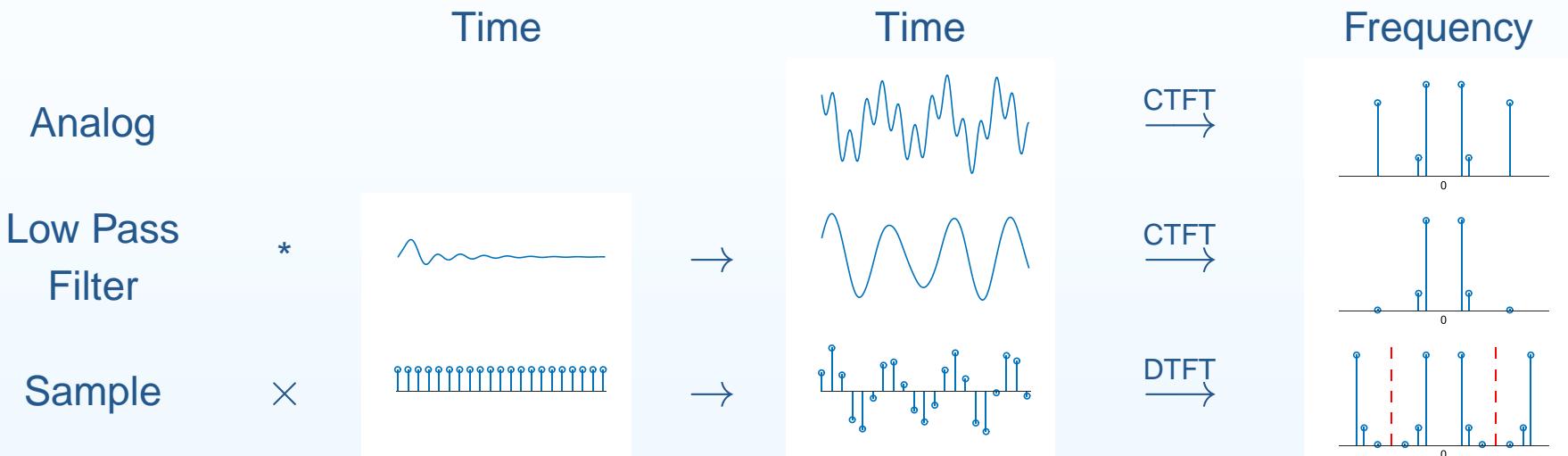
Sampling Process



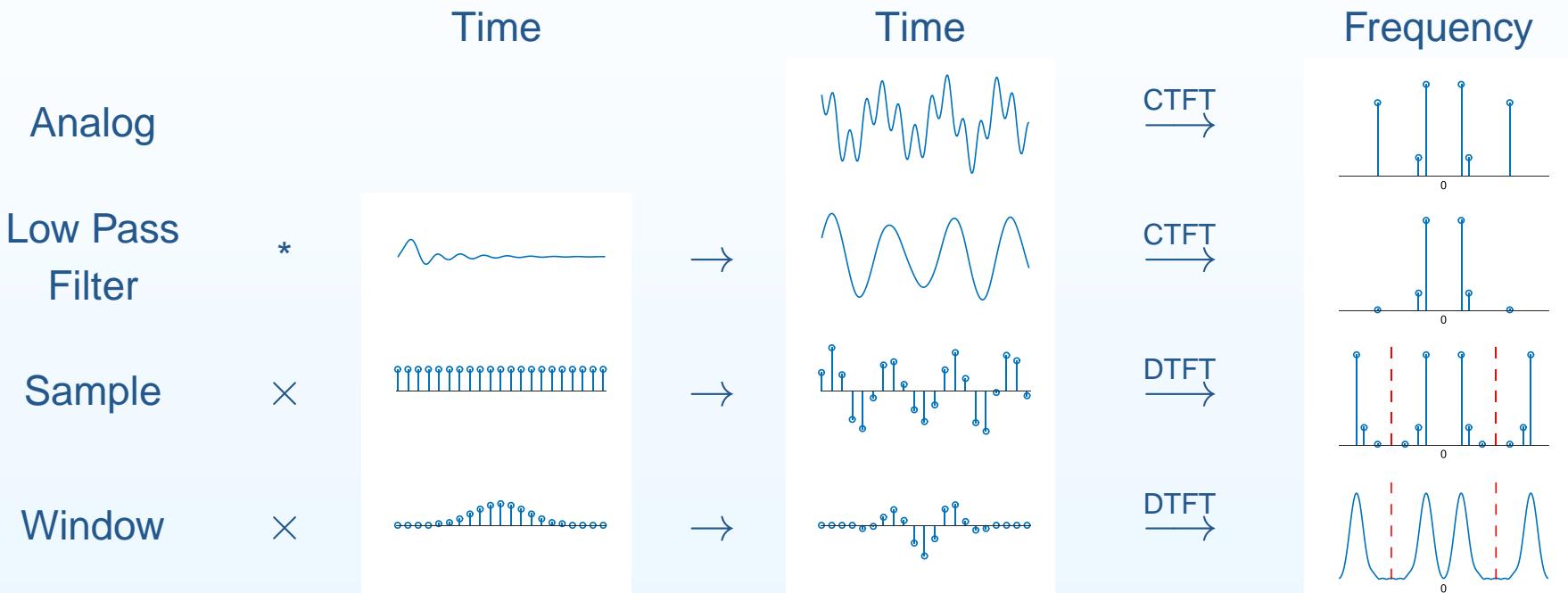
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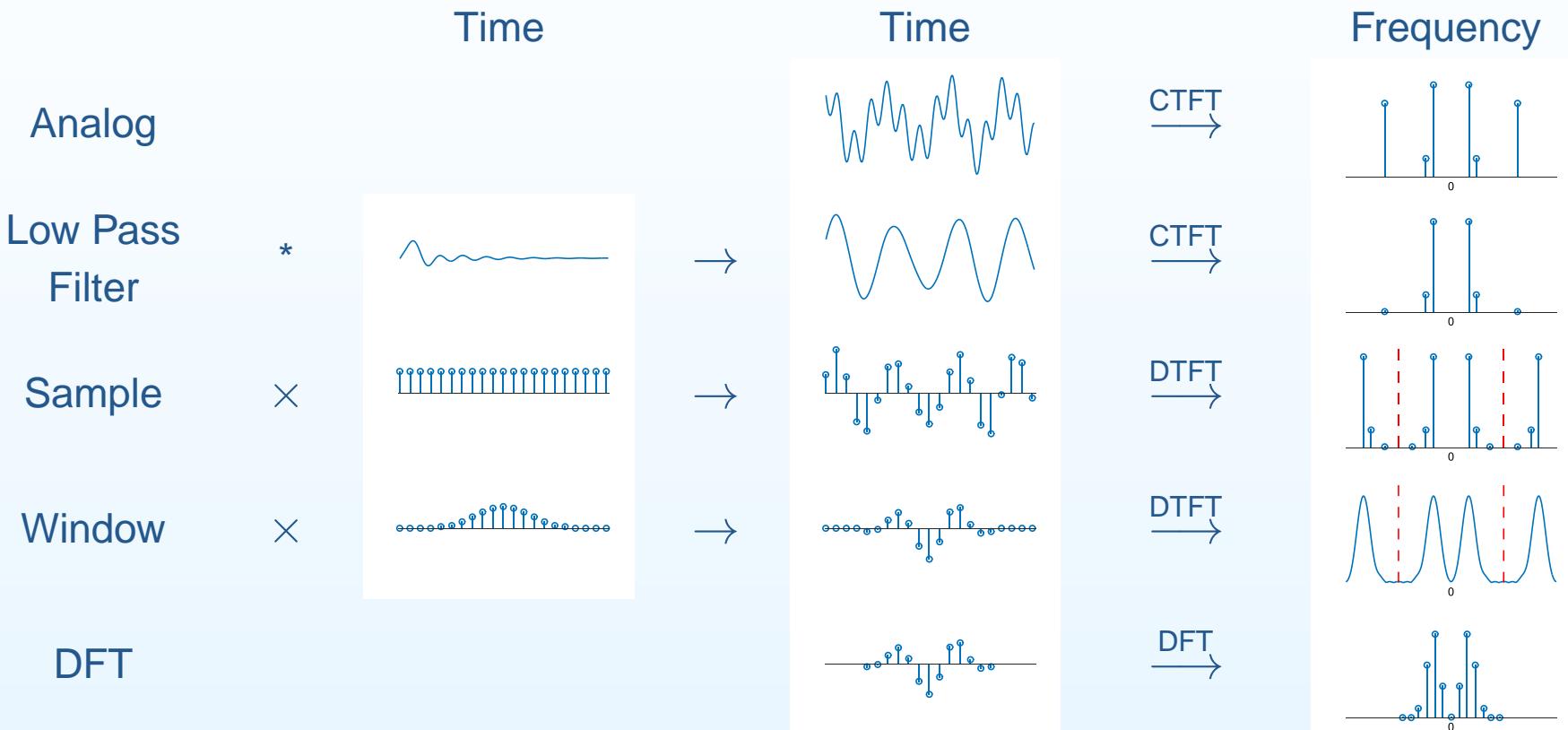
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Sampling Process



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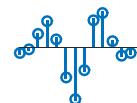
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Zero-Padding

Zero padding means added extra zeros onto the end of $x[n]$ before performing the DFT.

Windowed Signal

Time $x[n]$



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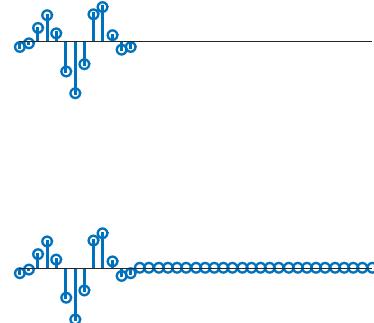
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Zero-Padding

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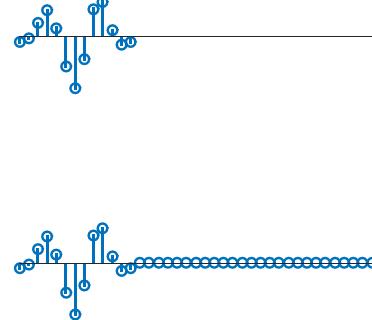
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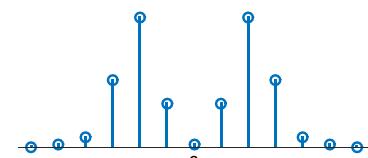
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Frequency $|X[k]|$



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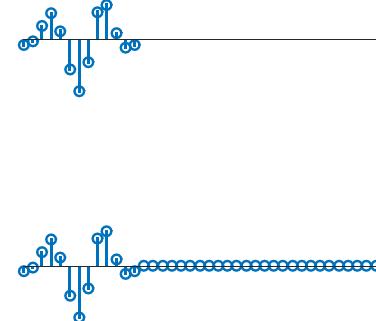
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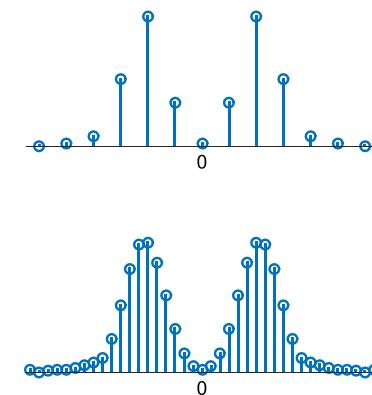
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Zero-Padding

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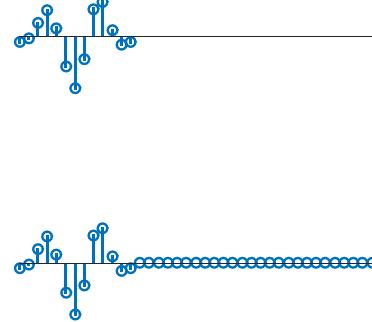
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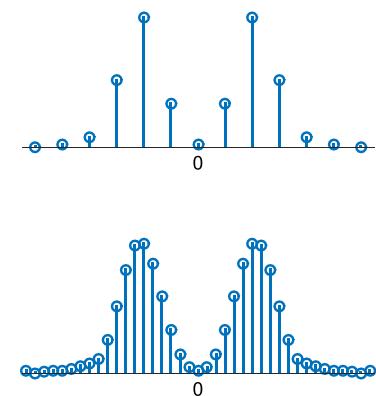
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Frequency $|X[k]|$



- Zero-padding causes the DFT to evaluate the DTFT at more values of ω_k . Denser frequency samples.

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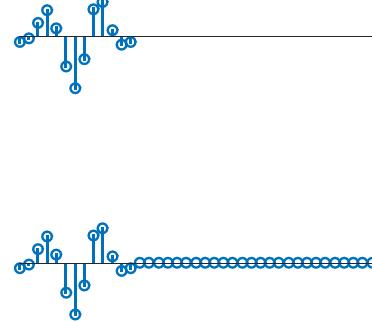
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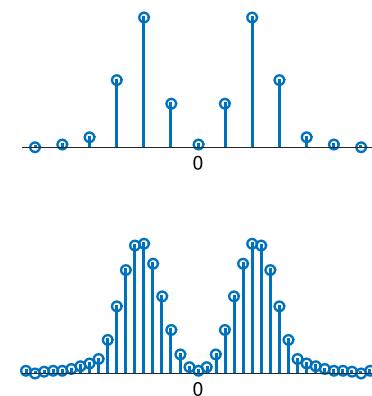
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With zero-padding

Time $x[n]$



Frequency $|X[k]|$



- Zero-padding causes the DFT to evaluate the DTFT at more values of ω_k . Denser frequency samples.
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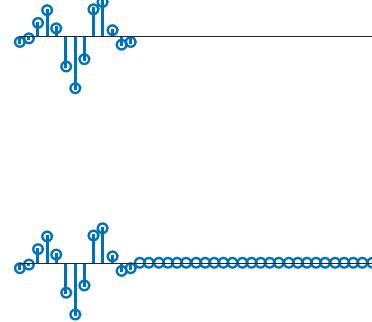
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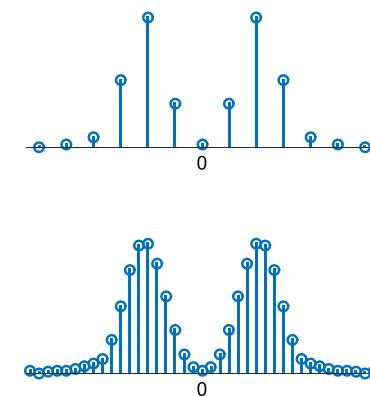
Windowed Signal

With zero-padding

Time $x[n]$



Frequency $|X[k]|$



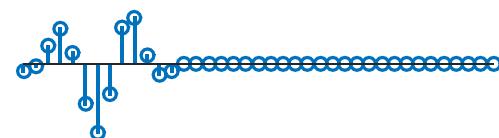
- Zero-padding causes the DFT to evaluate the DTFT at more values of ω_k . Denser frequency samples.
- Width of the peaks remains constant: determined by the length and shape of the window.
- Smoother graph but increased frequency resolution is an illusion.

Phase Unwrapping

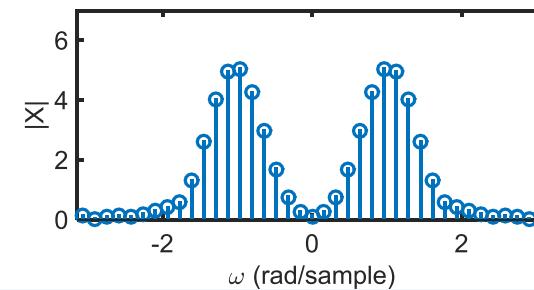
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Phase of a DTFT is only defined to within an integer multiple of 2π .



$x[n]$



$|X[k]|$

Phase Unwrapping

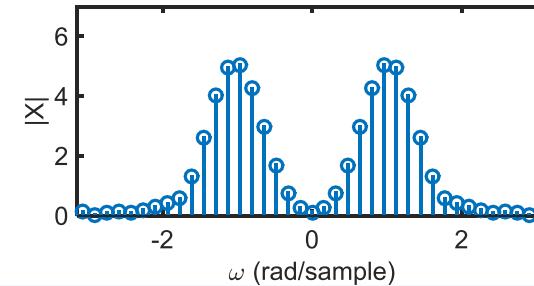
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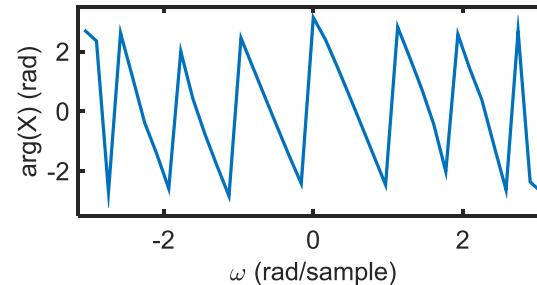
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$x[n]$



$|X[k]|$



$\angle X[k]$

Phase Unwrapping

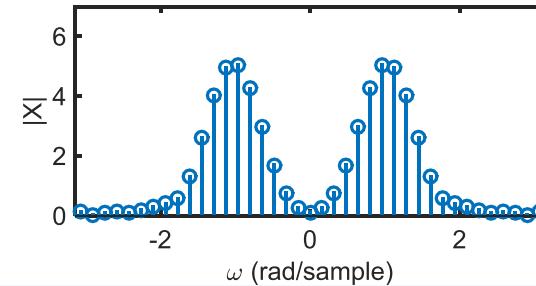
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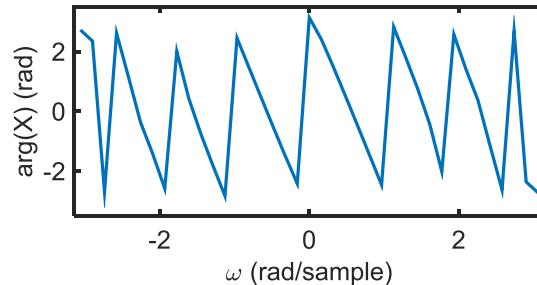
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$x[n]$



$|X[k]|$



$\angle X[k]$

Phase unwrapping adds multiples of 2π onto each $\angle X[k]$ to make the phase as continuous as possible.

Phase Unwrapping

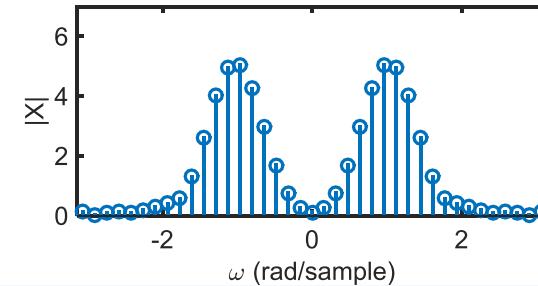
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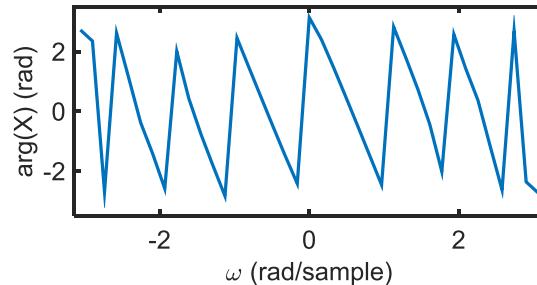
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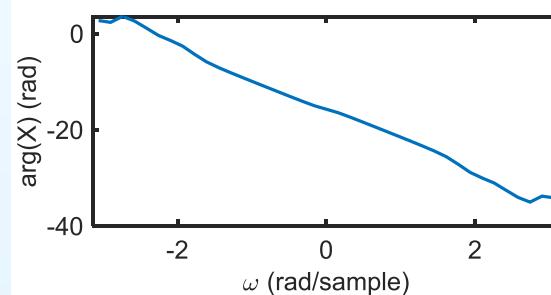
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$\angle X[k]$ unwrapped

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Assume $\int |x(t)|^2 dt = 1 \Rightarrow \int |X(j\omega)|^2 d\omega = 2\pi$ [Parseval]

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Now $\int tx \frac{dx}{dt} dt = \frac{1}{2}tx^2(t) \Big|_{t=-\infty}^{\infty} - \int \frac{1}{2}x^2 dt = 0 - \frac{1}{2}$ [by parts]

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So $\frac{1}{4} = \left| \int tx \frac{dx}{dt} dt \right|^2 \leq \left(\int t^2 x^2 dt \right) \left(\int \left| \frac{dx}{dt} \right|^2 dt \right)$ [Schwartz]

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$$\text{CTFT uncertainty principle: } \left(\frac{\int t^2 |x(t)|^2 dt}{\int |x(t)|^2 dt} \right)^{\frac{1}{2}} \left(\frac{\int \omega^2 |X(j\omega)|^2 d\omega}{\int |X(j\omega)|^2 d\omega} \right)^{\frac{1}{2}} \geq \frac{1}{2}$$

The first term measures the “width” of $x(t)$ around $t = 0$.

It is like σ if $|x(t)|^2$ was a zero-mean probability distribution.

The second term is similarly the “width” of $X(j\omega)$ in frequency.

A signal **cannot be concentrated in both time and frequency.**

Proof Outline:

Assume $\int |x(t)|^2 dt = 1 \Rightarrow \int |X(j\omega)|^2 d\omega = 2\pi$ [Parseval]

Set $v(t) = \frac{dx}{dt} \Rightarrow V(j\omega) = j\omega X(j\omega)$ [by parts]

Now $\int tx \frac{dx}{dt} dt = \frac{1}{2}tx^2(t) \Big|_{t=-\infty}^{\infty} - \int \frac{1}{2}x^2 dt = 0 - \frac{1}{2}$ [by parts]

So $\frac{1}{4} = \left| \int tx \frac{dx}{dt} dt \right|^2 \leq \left(\int t^2 x^2 dt \right) \left(\int \left| \frac{dx}{dt} \right|^2 dt \right)$ [Schwartz]
 $= \left(\int t^2 x^2 dt \right) \left(\int |v(t)|^2 dt \right)$

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No exact equivalent for DTFT/DFT but a similar effect is true

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- Three types: CTFT, DTFT, DFT

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- Three types: CTFT, DTFT, DFT
 - $\text{DTFT} = \text{CTFT of continuous signal} \times \text{impulse train}$

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 - $\frac{1}{2\pi}$ for CTFT and DTFT or $\frac{1}{N}$ for DFT
 - e.g. Inverse transform, Parseval, frequency domain convolution

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For further details see Mitra: 3 & 5.

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fft, ifft	DFT with optional zero-padding
fftshift	swap the two halves of a vector
conv	convolution or polynomial multiplication (not circular)
$x[n] \circledast y[n]$	<code>real(ifft(fft(x).*fft(y)))</code>
unwrap	remove 2π jumps from phase spectrum